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## On the Rectilinear Motion of Three Bodies Mutually Attracting Each Other

Sylvio R. Bistafa  
University of Sao Paulo, sbistafa@usp.br

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# On the rectilinear motion of three bodies mutually attracting each other

## E327 -- *De motu rectilineo trium corporum se mutuo attrahentium*

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Translated and Annotated

by

Sylvio R. Bistafa\*

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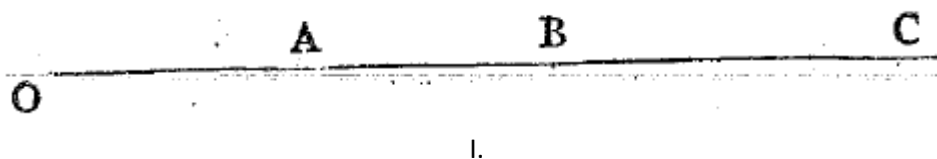
Foreword

This is an annotated translation from Latin of E327 -- *De motu rectilineo trium corporum se mutuo attrahentium* ("On the rectilinear motion of three bodies mutually attracting each other"). In this publication, Euler considers three bodies lying on a straight line, which are attracted to each other by central forces inversely proportional to the square of their separation distance (inverse-square law). Here Euler finds that the parameter that controls the relative distances among the bodies is given by a quintic function.

On the rectilinear motion of three bodies mutually attracting each other

Author

L. Euler



Let  $A, B, C$  be the masses of three bodies such that their distances to a fixed point  $O$  at a given instant of time  $t$  is given by

$$OA = x, OB = y \text{ and } OC = z$$

where, in fact, it is assumed that  $y > x$  and  $z > y$ . Hence, the principles of motion give these three equations:

$$I. \quad \frac{ddx}{dt^2} = \frac{B}{(y-x)^2} + \frac{C}{(z-x)^2} \quad [1]$$

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\* Corresponding address: sbistafa@usp.br

$$II. \quad \frac{ddy}{dt^2} = \frac{-A}{(y-x)^2} + \frac{C}{(z-y)^2} \quad [2]$$

$$III. \quad \frac{ddz}{dt^2} = \frac{-A}{(z-x)^2} - \frac{B}{(z-y)^2} \quad [3]$$

whence two integrable equations are easily derived: the first [integral]

$$Adx + Bdy + Cdz = Edt \quad [4]$$

which upon integration results in

$$Ax + By + Cz = Et + F;^1 \quad [5]$$

and the second [integral]

$$\frac{Adx^2 + Bdy^2 + Cdz^2}{dt^2} = G + \frac{2AB}{y-x} + \frac{2AC}{z-x} + \frac{2BC}{z-y}.^2 \quad [6]$$

Whence, because we lack a third integral equation, very little is possible to conclude about the movement.

2. Let us set  $x = y - p$  and  $z = y + q$ , so that  $p$  and  $q$  are positive quantities; and the first integral [Eq.5] gives:

$$(A + B + C)y - Ap + Cq = Et + F \quad [7]$$

and thus

$$y = \frac{Ap - Cq + Et + F}{A + B + C}; \quad dy = \frac{Adp - Cdq + Edt}{A + B + C} \quad [8 \text{ a, b}]$$

$$x = \frac{-(B + C)p - Cq + Et + F}{A + B + C}; \quad dx = \frac{-(B + C)dp - Cdq + Edt}{A + B + C} \quad [9 \text{ a, b}]$$

$$z = \frac{Ap + (A + B)q + Et + F}{A + B + C}; \quad dz = \frac{Adp + (A + B)dq + Edt}{A + B + C}. \quad [10 \text{ a, b}]$$

Whence, the second integral [Eq. 6] assumes the following form:

$$\begin{aligned} & \frac{A(B + C)dp^2 + C(A + B)dq^2 + 2ACdpdq + EEdt^2}{(A + B + C)dt^2} \\ & = G + \frac{2AB}{p} + \frac{2AC}{p + q} + \frac{2BC}{q}, \end{aligned} \quad [11]$$

<sup>1</sup> Upon the multiplication of I.[Eq. 1] by A; II. [Eq. 2] by B; III. [Eq. 3] by C, and taking the sum over the three equations yields  $Adx + Bdy + Cdz = Edt$ , which upon integration gives  $Ax + By + Cz = Et + F$ , where E and F are constants of integration.

<sup>2</sup> From I.  $\frac{Adx}{dt} = \int \frac{AB}{(y-x)^2} dt + \int \frac{AC}{(z-x)^2} dt$ , which upon multiplication by  $\frac{dx}{dt}$  gives  $\frac{Adx^2}{dt^2} = \int \frac{AB}{(y-x)^2} dx + \int \frac{AC}{(z-x)^2} dx$ , and then  $\frac{Adx^2}{dt^2} = \frac{AB}{y-x} + \frac{AC}{z-x}$ . And similarly for II. and III., yielding:  $\frac{Bdy^2}{dt^2} = \frac{AB}{y-x} + \frac{BC}{z-y}$ , and  $\frac{Cdz^2}{dt^2} = \frac{AC}{z-x} + \frac{BC}{z-y}$ . Taking the sum over the three equations and introducing the constant of integration G yields  $\frac{Adx^2+Bdy^2+Cdz^2}{dt^2} = G + \frac{2AB}{y-x} + \frac{2AC}{z-x} + \frac{2BC}{z-y}$ .

whence arises one integral equation<sup>3</sup>

$$\frac{B(Adp^2 + Cdq^2) + AC(dp + dq)^2}{(A + B + C)dt^2} = G + \frac{2AB}{p} + \frac{2AC}{p + q} + \frac{2BC}{q}, \quad [12]$$

where the last term  $EE$  is included into  $G$ .

3. Let us, indeed, make the same substitutions into the first second order differential [difference-of- differentials] equations [Eqs. 1 and 3], which now result in two [equations]:

$$\frac{-(B + C)ddp - Cddq}{(A + B + C)dt^2} = \frac{B}{pp} + \frac{C}{(p + q)^2} \quad [13]$$

$$\frac{Addp + (A + B)ddq}{(A + B + C)dt^2} = \frac{-A}{(p + q)^2} - \frac{B}{qq}, \quad [14]$$

which [by subtracting Eq. 13 from Eq. 14] results in

$$\frac{ddp + ddq}{dt^2} = \frac{-A - C}{(p + q)^2} - \frac{B}{pp} - \frac{B}{qq}. \quad [15]$$

And then, each element  $ddp$  and  $ddq$  can be expressed separately in the following way

$$1^{\circ}. \quad \frac{ddp}{dt^2} = \frac{-A - B}{pp} - \frac{C}{(p + q)^2} + \frac{C}{qq} \quad [16]$$

$$2^{\circ}. \quad \frac{ddq}{dt^2} = \frac{A}{pp} - \frac{A}{(p + q)^2} - \frac{B + C}{qq}. \quad [17]^4$$

4. Since the solution has been reduced to two second-order differential equations involving  $p$ ,  $q$  and  $t$  we should expect that significant advantage is to be obtained, if it were possible to reduce these equation to two others of first order only. This is a unique technique that I have discovered which can be applied in the following manner. I put  $q = pu$ , and the two differential equations [Eqs. 16 and 17] are represented as:

$$d\left(\frac{dp}{dt}\right) = \frac{dt}{pp} \left(-A - B - \frac{C}{(u + 1)^2} + \frac{C}{uu}\right) \quad [18]$$

$$d\left(\frac{udp + pdu}{dt}\right) = \frac{dt}{pp} \left(A - \frac{A}{(u + 1)^2} - \frac{B + C}{uu}\right). \quad [19]$$

Now the trick consists of the following substitution  $\frac{dp}{dt} = \frac{r}{\sqrt{p}}$  and  $\frac{dq}{dt} = \frac{udp + pdu}{dt} = \frac{s}{\sqrt{p}}$ ; indeed it will become clear that by means of these substitutions, the two variables  $p$  and  $t$  can be eliminated from the calculations, such that only these three [variables]  $r$ ,  $s$  and  $u$  are to be determined by their first differentials. Then, in

<sup>3</sup> In the original paper (*Novi Commentarii academiae scientiarum Petropolitanae* 11, 1767, pp. 144 – 151), these three last lines have been misplaced at the end of § 3.

<sup>4</sup> There is a sign misprint in the original article. This error occurred elsewhere, in equations [22], [24], [27], [28], [29].

particular, the integral equation found above [Eq. 12] assumes a finite form which reads

$$\frac{B(Arr + Css) + AC(r + s)^2}{A + B + C} = Gp + 2AB + \frac{2AC}{u + 1} + \frac{2BC}{u}, \quad [20]$$

which will be able to produce remarkable applications.

5. Since  $\frac{dp}{dt} = \frac{r}{\sqrt{p}}$ , then  $dt = \frac{dp\sqrt{p}}{r}$ , whence our second order differential [difference-of-differentials] equations [Eqs. 18 and 19] give

$$\frac{dr}{\sqrt{p}} - \frac{rdp}{2p\sqrt{p}} = \frac{dp}{pr\sqrt{p}} \left( -A - B - \frac{C}{(u + 1)^2} + \frac{C}{uu} \right) \quad [21]$$

$$\frac{ds}{\sqrt{p}} - \frac{sdp}{2p\sqrt{p}} = \frac{dp}{pr\sqrt{p}} \left( A - \frac{A}{(u + 1)^2} - \frac{B + C}{uu} \right). \quad [22]$$

or:

$$dr = \frac{rdp}{2p} + \frac{dp}{pr} \left( -A - B - \frac{C}{(u + 1)^2} + \frac{C}{uu} \right) \quad [23]$$

$$ds = \frac{sdp}{2p} + \frac{dp}{pr} \left( A - \frac{A}{(u + 1)^2} - \frac{B + C}{uu} \right). \quad [24]$$

Moreover, in particular, it will be considered that

$$udp + pdu = \frac{sdt}{\sqrt{p}} = \frac{sdp}{r}, \quad [25]$$

and thus  $\frac{dp}{p} = \frac{rdu}{s - ru}$  which when substituted [into Eqs. 23 and 24] gives

$$dr(s - ru) = \frac{1}{2}rrdu + du \left( -A - B - \frac{C}{(u + 1)^2} + \frac{C}{uu} \right) \quad [26]$$

$$ds(s - ru) = \frac{1}{2}rsdu + du \left( A - \frac{A}{(u + 1)^2} - \frac{B + C}{uu} \right), \quad [27]$$

which when combined give:

$$\begin{aligned} \frac{1}{2}r(rds - sdr) + ds \left( -A - B - \frac{C}{(u + 1)^2} + \frac{C}{uu} \right) - dr \left( A - \frac{A}{(u + 1)^2} - \frac{B + C}{uu} \right) \\ = 0. \end{aligned} \quad [28]$$

6. We see that we have two first-order differential equations [Eqs. 26 and 27] among three variables  $r$ ,  $s$  and  $u$ , whence if it were possible to determine  $r$  and  $s$  in terms of  $u$ , and then one would have the complete solution of the problem. Thence, in fact,  $p$  would become known from the formula  $\frac{dp}{p} = \frac{rdu}{s - ru}$ , and hence furthermore  $q = pu$ . Thereafter, the particular time  $t$  would be given from the equation  $dt = \frac{dp\sqrt{p}}{r} = \frac{pdu\sqrt{p}}{s - ru}$ ; and finally, for a given time  $t$ , the distances  $x, y, z$  would be obtained as given in § 2.

7. Since the two differential equations [Eqs. 26 and 27] found are

$$dr(s - ru) = \frac{1}{2}rrdu + du\left(-A - B - \frac{C}{(u + 1)^2} + \frac{C}{uu}\right)$$

$$ds(s - ru) = \frac{1}{2}rsdu + du\left(A - \frac{A}{(u + 1)^2} - \frac{B + C}{uu}\right),$$

then, it is clear that both are satisfied by taking the quantity  $u$  constant and  $s - ru = 0$ , whence a particular solution can be obtained by setting  $u = \alpha$  and  $s = \alpha r$

$$-(A + B) \alpha - \frac{C \alpha}{(\alpha + 1)^2} + \frac{C}{\alpha} = A - \frac{A}{(\alpha + 1)^2} - \frac{B + C}{\alpha \alpha}, \quad [29]$$

or

$$0 = A\left(\alpha + 1 - \frac{1}{(\alpha + 1)^2}\right) + B\left(\alpha - \frac{1}{\alpha \alpha}\right) + C\left(\frac{\alpha}{(\alpha + 1)^2} - \frac{1}{\alpha} - \frac{1}{\alpha \alpha}\right), \quad [30]$$

or

$$0 = A \frac{[(\alpha + 1)^3 - 1]}{(\alpha + 1)^2} + \frac{B(\alpha^3 - 1)}{\alpha \alpha} + \frac{C[\alpha^3 - (\alpha + 1)^3]}{\alpha \alpha (\alpha + 1)^2}; \quad [31]$$

hence,

$$C(1 + 3 \alpha + 3 \alpha \alpha) = A \alpha^3 (\alpha \alpha + 3 \alpha + 3) + B(\alpha + 1)^2(\alpha^3 - 1). \quad [32]$$

Thus, it is possible to determine the quantity  $\alpha$  from this equation of the fifth degree:

$$(A + B) \alpha^5 + (3A + 2B) \alpha^4 + (3A + B) \alpha^3 - (B + 3C) \alpha^2 - (2B + 3C) \alpha - B - C = 0. \quad [33]$$

Consequently, the relation between  $r$  and  $p$  [Eq. 21] is given by the following equation

$$dr = \frac{rdp}{2p} + \frac{dp}{pr}\left(-A - B - \frac{C}{(\alpha + 1)^2} + \frac{C}{\alpha \alpha}\right), \quad [34]$$

or by putting  $A + B + \frac{C}{(\alpha + 1)^2} - \frac{C}{\alpha \alpha} = \frac{1}{2}D$ , [then, from Eq. 34 we have that]

$$2dr = \frac{dp}{p}\left(r - \frac{D}{r}\right) \text{ or } \frac{dp}{p} = \frac{2rdr}{rr - D}, \quad [35 \text{ a, b}]$$

which [upon integration of Eq. 35b] gives

$$p = \beta(rr - D), \quad [36]$$

[where  $\beta$  is a constant of integration], and then [since  $q = pu = p \alpha$ ]

$$q = \alpha \beta(rr - D), \quad [37]$$

and [since]  $dt = \frac{dp\sqrt{p}}{r}$  or  $dt = 2\beta dr\sqrt{\beta(rr - D)}$ , hence

$$t = \beta r \sqrt{\beta(rr - D)} - \beta^2 D \int \frac{dr}{\sqrt{\beta(rr - D)}}. \quad [38]$$

8. This particular case, in which the solution succeeds, deserves to be unfolded more carefully. Firstly, I observe that the value of  $\alpha$ , since it is obtained from an

equation of the fifth degree, is unique and always a positive real quantity, because there is only one sign variation<sup>5</sup>, and then, of course, there is no reason for any ambiguity; consequently,  $\alpha$  is determined by the masses  $A, B, C$ . Once the number  $\alpha$  having been found, we get the quantity  $D = 2(A + B) - \frac{2C(2\alpha+1)}{\alpha\alpha(\alpha+1)^2}$ , where it should be observed that the quantity  $D$  can never vanish. In fact, if  $D = 0$ , then  $B = \frac{C(2\alpha+1)}{\alpha\alpha(\alpha+1)^2} - A$ , which, when substituted [into Eq. 32] would give:

$$\begin{aligned}
 & C(1 + 3\alpha + 3\alpha\alpha) \\
 &= A\alpha^3(\alpha\alpha + 3\alpha + 3) + \frac{C(2\alpha + 1)(\alpha^3 - 1)}{\alpha\alpha} \\
 &- A(\alpha + 1)^2(\alpha^3 - 1), \tag{39}
 \end{aligned}$$

or

$$\frac{C}{\alpha\alpha}(\alpha^4 + 2\alpha^3 + \alpha\alpha + 2\alpha + 1) = A(\alpha^4 + 2\alpha^3 + \alpha\alpha + 2\alpha + 1), \tag{40}$$

and, therefore,  $C = A\alpha\alpha$  and  $B = \frac{A(2\alpha+1)}{(\alpha+1)^2} - A = \frac{-A\alpha\alpha}{(\alpha+1)^2}$ , and then,  $B$  would be a negative mass, which is absurd. Even less possible is that the quantity  $D$  could be negative. In fact, assuming [that  $\frac{D}{2} = -\Delta$ , where  $\Delta$  is a positive quantity ]:

$$B = \frac{C(2\alpha + 1)}{\alpha\alpha(\alpha + 1)^2} - A - \Delta, \tag{41}$$

it would then give [when substituted into Eq. 32]:

$$\frac{C}{\alpha\alpha} = A - \frac{\Delta(\alpha + 1)^2(\alpha^3 - 1)}{\alpha^4 + 2\alpha^3 + \alpha\alpha + 2\alpha + 1}, \tag{42}$$

hence [by isolating  $A$  in the first term of Eq. 42, and substituting the resulting expression into Eq. 41]

$$B = \frac{C(2\alpha + 1)}{\alpha\alpha(\alpha + 1)^2} - \frac{C}{\alpha\alpha} - \frac{\Delta(\alpha^5 + 3\alpha^4 + 3\alpha^3)}{\alpha^4 + 2\alpha^3 + \alpha\alpha + 2\alpha + 1}, \tag{43}$$

and then,  $B$  would be a much more negative quantity, since it is necessary that the value of  $\alpha$  itself be positive.

9. Then, since it is necessary that the quantity  $D$  be positive, it can be assumed that  $D = aa$ , and if also the number  $\alpha$  is considered as given, therefore the masses of the three bodies will be obtained as [from Eq. 43, with  $\Delta = -\frac{aa}{2}$ ]:

$$B = \frac{(\alpha^5 + 3\alpha^4 + 3\alpha^3)aa}{2(\alpha^4 + 2\alpha^3 + \alpha\alpha + 2\alpha + 1)} - \frac{C}{(\alpha + 1)^2}; \tag{44}$$

and [from Eq. 42, with  $\Delta = -\frac{aa}{2}$ ]:

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<sup>5</sup> Descartes' Rule of Signs states that if the terms of a single-variable polynomial with real coefficients are ordered by descending variable exponent, then the number of positive roots of the polynomial is either equal to the number of sign differences between consecutive nonzero coefficients, or is less than it by an even number.

$$A = \frac{C}{\alpha\alpha} - \frac{(\alpha + 1)^2(\alpha^3 - 1)aa}{2(\alpha^4 + 2\alpha^3 + \alpha\alpha + 2\alpha + 1)}, \quad [45]$$

from which it is necessary that the quantity  $\frac{2C(\alpha^4 + 2\alpha^3 + \alpha\alpha + 2\alpha + 1)}{\alpha\alpha(\alpha + 1)^2aa}$  be bound between the limits  $(\alpha + 1)^3 - 1$  and  $\alpha^3 - 1$ . Then, once the quantity  $aa$  with the number  $\alpha$  are introduced into the calculations, two cases should be examined, according to the sign of the quantity  $\beta$  (positive or negative), which we shall examine separately.

Case I.

10. First let  $\beta = nn$ , and then  $p = nn(rr - aa)$  [from Eq. 36] and  $q = \alpha nn(rr - aa)$  [from Eq. 37], and we can set the constants  $E$  and  $F$  to zero, the locations of the three bodies  $A, B, C$ , whose center of gravity is now located in  $O$ , are defined by  $r$  such that:

$$x = OA = \frac{-nn(rr - aa)}{A + B + C} (B + C + C\alpha), \quad [46]$$

$$y = OB = \frac{nn(rr - aa)}{A + B + C} (A - C\alpha), \quad [47]$$

$$z = OC = \frac{nn(rr - aa)}{A + B + C} (A + (A + B)\alpha). \quad [48]$$

Yet, the relation between  $r$  and the time  $t$  is [Eq. 38]

$$t = n^3r\sqrt{rr - aa} - n^3aa \int \frac{dr}{\sqrt{rr - aa}}, \quad [49]$$

or

$$t = n^3r\sqrt{rr - aa} - n^3aa \ln \frac{r + \sqrt{rr - aa}}{\Delta}. \quad [50]$$

Assuming that the constant  $\Delta = a$ , then, for the time  $t = 0$ ,  $r = a$ , meaning that all bodies are joined in the center of gravity  $[O]$ , whence they will be driven out with almost infinite velocities, and then, these [distances] are similar to each other as the quantities of:  $-B - C - C\alpha$ ,  $A - C\alpha$ ,  $A + (A + B)\alpha$ ; also, with the passage of time  $t$  the quantity  $r$  increases even more; however, for any other time, the velocity of each of the bodies becomes known from the formula  $\frac{dt}{dr} = 2n^3\sqrt{rr - aa}$ . However, it should be noted that the inter body distances preserve the same proportion.

Case II.

11. Now let  $\beta = -nn$ , and then  $p = nn(aa - rr)$  [from Eq. 36] and  $q = \alpha nn(aa - rr)$  [from Eq. 37], and the locations of the bodies are defined by  $r$  from:

$$x = OA = \frac{-nn(aa - rr)}{A + B + C} (B + (1 + \alpha)C), \quad [51]$$

$$y = OB = \frac{nn(aa - rr)}{A + B + C} (A - C\alpha), \quad [52]$$



$$z = OC = \frac{nn(aa - rr)}{A + B + C} (A(\alpha + 1) + B \alpha). \tag{53}$$

On the other hand, for the time  $t$  to be obtained,  $dt = 2n^3 dr \sqrt{(aa - rrr)}$ ,

or

$$t = n^3 r \sqrt{aa - rr} + n^3 aa \int \frac{dr}{\sqrt{aa - rr}}, \tag{54}$$

hence,

$$t = n^3 r \sqrt{aa - rr} + n^3 aa \sin^{-1} \left( \frac{r}{a} \right). \tag{55}$$

But, by putting  $\sin^{-1} \left( \frac{r}{a} \right) = \phi$ , such that  $r = a \sin \phi$ , then  $t = n^3 aa (\phi + \sin \phi \cos \phi)$ , and at any time, the relative distances are proportional to  $\cos^2 \phi$ .<sup>6</sup> Whence, if at the beginning when  $t = 0$ , also  $\phi = 0$ , thus  $r = 0$ , and  $\frac{dt}{dr} = 2n^3 a$ , then the distances will be:

$$x = OA = \frac{-nnaa}{A + B + C} (B + (1 + \alpha)C), \tag{56}$$

$$y = OB = \frac{nnaa}{A + B + C} (A - C \alpha), \tag{57}$$

$$z = OC = \frac{nnaa}{A + B + C} (A(\alpha + 1) + B \alpha), \tag{58}$$

and in that place the bodies are at rest. On the other hand, once it has been assumed that  $\phi = 90^\circ$ , or after the time  $t = n^3 aa \cdot 90^\circ$  has elapsed, the bodies approach the center of gravity with infinite velocity.<sup>7</sup>

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### An Application by the Translator

For  $A = B = C = 1$  unit, the fifth order polynomial [Eq.33] reduces to

$$2 \alpha^5 + 5 \alpha^4 + 4 \alpha^3 - 4 \alpha^2 - 5 \alpha - 2 = 0,$$

which, by inspection, gives  $\alpha = 1$  as the single positive root.

For Case I, and from Eqs. 46-48, the positions of these masses are given by  $x = -nn(rr - aa)$ ,  $y = 0$ , and  $z = nn(rr - aa)$ , which show that one of the masses occupies the center of gravity, whereas the two other masses occupy symmetrical positions in relation to the center of gravity, and that these positions depend on  $(rr - aa)$ .

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<sup>6</sup> According to Eqs. 51-53, the distances are proportional to  $(a^2 - r^2) = \left( a^2 - a^2 \frac{r^2}{a^2} \right) = a^2(1 - \sin^2 \phi) = a^2 \cos^2 \phi$ .

<sup>7</sup> Since  $dt = 2nndr \sqrt{nn(aa - rr)}$ , or  $\frac{dr}{dt} = -\frac{1}{2n^3 \sqrt{(aa - rr)}}$ , which for  $\phi = 90^\circ$  ( $r = a$ ) gives  $\frac{dr}{dt} = -\infty$ .

For Case II, and when  $t = 0$ , the motion begins in a symmetrical configuration of the masses in relation to the center of gravity, however, according to Eqs. 56-58, the positions of these masses are now given by  $x = -nnaa$ ,  $y = 0$ , and  $z = nnaa$ . For any other instant of time, the positions of the masses will be given by  $x = -nn(aa - rr)$ ,  $y = 0$ , and  $z = nn(aa - rr)$ , which show that the symmetrical configuration is preserved, however, the positions of the masses now depend on  $(aa - rr)$  instead.

The conclusion is that for these particular cases, the configuration of the masses during the motion is such that one of the masses occupies the center of gravity, with the two other masses remaining on the same straight line, and moving symmetrically around the center of gravity of the system. As shown earlier, the quintic equation has a unique positive root, even when the masses of the bodies are not equal. In this case, this could represent a hypothetical situation of Sun, Earth and Moon in perpetual alignment in syzygy, for which the parameter that controls the relative distances among the planets was found to be given by a quintic function.